BACKGROUND QUANTUM GEOMETRY FLAT BANDS PLASMONS FUTURE RESEARCH 00 000 000 000 000 00

Manifestations of Quantum Geometric Effects

Lydia England

Advisor: Dr. Bruno Uchoa

University of Oklahoma REU 2024



July 31, 2024

References

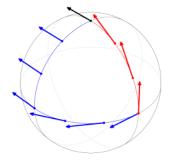
PARALLEL TRANSPORT AND ANHOLONOMY

FLAT BANDS

Two vectors are displaced along paths drawn on a sphere & maintain a constant angle with the surface *locally*.

BACKGROUND

- The vectors in the figure "rotate" with respect to *each other*, but are parallel to the sphere during transport.
- This global rotation without any accompanying local rotation is a result of the intrinsic curvature of the sphere.



FUTURE RESEARCH

Figure: Parallel transport

Dac

イロト イ理ト イヨト イヨト



- 1931 Dirac writes about a path-dependent non-integrable phase in a paper titled "Quantised Signularities in the Electromagnetic Field" [1].
 - 1980 Provost & Vallee introduce a Riemannian metric structure on the manifold of quantum states [2].
 - 1984 M. Berry proposes the quantum geometric phase in a seminal paper [3].



QUANTUM GEOMETRIC TENSOR

OUANTUM GEOMETRY

- The state of a quantum system is described by a ray (i.e., equivalence class of vectors which differ only by a phase) in a complex (projective) Hilbert space, P(H).
- For $\mathcal{H} = \mathcal{H}_n$ finite-dimensional, the space $\mathbf{P}(\mathcal{H}_n)$ can equivalently be treated as a complex projective space: $\mathbf{P}(\mathcal{H}_n) = \mathbb{C}\mathbf{P}^{n-1}$ and structured as a smooth manifold.
 - CPⁿ⁻¹ is naturally endowed with the Fubini-Study metric³D_{1.2.3}
 this carries mutually compatible complex, Riemannian, and symplectic structures.

・ ロ ト ス 厚 ト ス 戸 ト ス 戸 ト

QUANTUM GEOMETRIC TENSOR

QUANTUM GEOMETRY

The naturally emerging metric structure is the Quantum Geometric Tensor (QGT).

4p4d

$$T_{ij}(u_{\lambda}) \equiv \left\langle \partial_{i} u_{\lambda} \mid \left(1 - \left| u_{\lambda} \right\rangle \left\langle u_{\lambda} \right| \right) \mid \partial_{j} u_{\lambda} \right\rangle$$

 $(\partial_i = \partial/\partial(k_i)$ with $i = x, y \& u_\lambda$ is a wave function parametrized by λ).

The QGT is simply the Fubini-Study metric on the $\mathbb{C}\mathbf{P}^{n-1}$ manifold (equivalently, the projective Hilbert space $\mathbf{P}(\mathcal{H}_n)$).

nac

QUANTUM GEOMETRY

Quantum states are represented by complex functions; hence, the QGT is complex.

4p4d

$$T_{ij}\equiv g_{ij}+rac{i}{2}\Omega_{ij}$$

• The real (symmetric) part, g_{ij} , is the quantum metric.

The imaginary (antisymmetric) part, Ω_{ij}, is the Berry curvature.

ELECTRONIC BAND STRUCTURE

BACKGROUND

The electronic band structure describes the allowed range of energy levels that electrons may occupy.

FLAT BANDS

FUTURE RESEARCH

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □

References

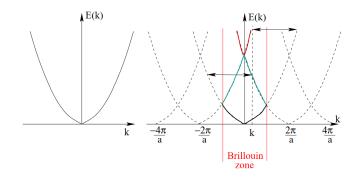


Figure: Free space eigenstates.

WHAT ARE FLAT ELECTRONIC BANDS?

FLAT BANDS

$$E(\mathbf{k}) = \frac{\hbar^2 k^2}{2m_*}$$

k wavevectorm_{*} effective mass

The energy range spanned by the entire band is called the bandwidth Δ . Flat bands are systems where the energy of single electrons energy *does not depend on momentum*.

• $m_* \to \infty$

- zero velocity
- electrons localized
- Fermi surface not well defined

(ロト (雪ト (ヨト (ヨ))

•
$$\Delta \rightarrow 0$$
.

WHY ARE FLAT ELECTRONIC BANDS INTERESTING?

FLAT BANDS

- In typical electronic band structures, we assume the electron interaction energy scale is *much smaller than* Δ .
 - ► Hence, electron-electron interactions are usually neglected.
- In a flat electronic band, $\Delta \rightarrow 0$.
 - Electron interactions can become the dominant interactions in the system.

・ロト ・ 母 ト ・ ヨ ト ・ ヨ ト

WHY ARE FLAT ELECTRONIC BANDS INTERESTING?

Quantum geometric effects become prominent in systems with flat electronic bands.

- Lower bound for stiffness in superconductors.
 - ► Transport despite no group velocity.
 - High superconducting critical temperature (twisted bilayer graphene).^S = 1

ロトス得トスラトスラト



- 1900 Paul Drude develops the Drude model of electrical conduction [4].
 - ► **1933:** Sommerfeld & Bethe incorporate quantum mechanics into the Drude model.

・ロト ・ 理 ト ・ ヨ ト

200

Early models treat interactions within a collection of electrons based on a free particle approximation.

However, this approximation mostly neglects the long-range correlations of electron positions which result from Coulomb interactions.

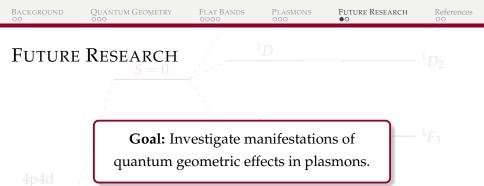


- 1952 David Bohm & David Pines develop a "collective" treatment of interactions in a collection of electrons,
 describing long-range Coulomb interactions of electrons in terms of collective fields which represent organized so-called "plasma" oscillations of the whole system [5–8].
 - 1956 David Pines coins the term "plasmon" in a 1956 review article [9].

 $\frac{B_{ACKGROUND}}{OOO} \underbrace{\begin{array}{c} QUANTUM GEOMETRY \\ OOO\end{array}}_{OOO} \underbrace{\begin{array}{c} FLAT BANDS \\ OOO\end{array}}_{OOO} \underbrace{\begin{array}{c} PLASMONS \\ OOO\end{array}}_{OO} \underbrace{\begin{array}{c} FUTURE RESEARCH \\ OOO\end{array}}_{OOO} \underbrace{\begin{array}{c} References \\ OOO\end{array}}_{OO} \underbrace{\begin{array}{c} ID \\ ID \\ D2 \end{array}}_{D2}$

More recently (2021), it has been demonstrated that there may be nontrivial internal structure within plasmon wave functions; moreover, that this may be intimately tied to their quantum geometry [10].

In [10], the authors considered the effects of the Berry curvature — the imaginary (antisymmetric) part of the QGT.



Challenge: finding flat band systems where plasmons occur.

200

In particular, how does the quantum metric manifest in plasmons?



References I

- [1] Paul Adrien Maurice Dirac. "Quantised Singularities in the Electromagnetic Field". In: Proceedings of The Royal Society A: Mathematical, Physical and Engineering Sciences 133 (1931), pp. 60–72. URL: https://api.semanticscholar.org/CorpusID:27548208.
- [2] J. P. Provost and G. Vallee. "Riemannian structure on manifolds of quantum states". In: Communications in Mathematical Physics 76.3 (1980), pp. 289–301.
- [3] Michael Victor Berry. "Quantal phase factors accompanying adiabatic changes". In: Proc.Roy.Soc.Lond. A 392.1802 (Mar. 1984), pp. 45–57. URL: https://www.jstor.org/stable/2397741.
- [4] P. Drude. "Zur Elektronentheorie der Metalle". In: Annalen der Physik 306.3 (Jan. 1900), pp. 566–613. ISSN: 1521-3889. DOI: 10.1002/andp.19003060312. URL: http://dx.doi.org/10.1002/andp.19003060312.
- [5] David Bohm and David Pines. "A Collective Description of Electron Interactions. I. Magnetic Interactions". In: Physical Review 82.5 (June 1951), pp. 625–634. ISSN: 0031-899X. DOI: 10.1103/physrev.82.625. URL: http://dx.doi.org/10.1103/PhysRev.82.625.
- [6] David Pines and David Bohm. "A Collective Description of Electron Interactions: II. Collective vs Individual Particle Aspects of the Interactions". In: *Physical Review* 85.2 (Jan. 1952), pp. 338–353. ISSN: 0031-899X. DOI: 10.1103/physrev.85.338. URL: http://dx.doi.org/10.1103/PhysRev.85.338.

References

REFERENCES II

- [7] David Bohm and David Pines. "A Collective Description of Electron Interactions: III. Coulomb Interactions in a Degenerate Electron Gas". In: *Physical Review* 92.3 (Nov. 1953), pp. 609–625. ISSN: 0031-899X. DOI: 10.1103/physrev.92.609. URL: http://dx.doi.org/10.1103/PhysRev.92.609.
- [8] David Pines. "A Collective Description of Electron Interactions: IV. Electron Interaction in Metals". In: Physical Review 92.3 (Nov. 1953), pp. 626–636. ISSN: 0031-899X. DOI: 10.1103/physrev.92.626. URL: http://dx.doi.org/10.1103/PhysRev.92.626.
- [9] David Pines. "Collective Energy Losses in Solids". In: Reviews of Modern Physics 28.3 (July 1956), pp. 184–198. ISSN: 0034-6861. DOI: 10.1103/revmodphys.28.184. URL: http://dx.doi.org/10.1103/RevModPhys.28.184.
- [10] Jinlyu Cao, H.A. Fertig, and Luis Brey. "Quantum Internal Structure of Plasmons". In: Physical Review Letters 127.19 (Nov. 2021). ISSN: 1079-7114. DOI: 10.1103/physrevlett.127.196403. URL: http://dx.doi.org/10.1103/PhysRevLett.127.196403.

イロト イ理ト イヨト イヨト

1

Dac